Expectation Values

Quantum Information and Computing Seminar

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Qubits

A **classical bit** is an abstraction of a classical binary event (such as electricity flowing or not flowing through a wire). We label one of the events as 0 and the other as 1.

A **qubit**, or quantum bit, is an abstraction of a quantum binary events (such as an electron being spin-up or spin-down). Therefore, they are represented as discrete quantum states as opposed to just the numbers 0 and 1. Mathematically they are

$$|0\rangle = \begin{pmatrix} 1\\ 0 \end{pmatrix}$$
$$|1\rangle = \begin{pmatrix} 0\\ 1 \end{pmatrix}$$

The **quantum state** of a qubit is represented by a complex-valued, 2×1 vector. It can be written as a linear combination of $|0\rangle$ and $|1\rangle$ as follows

$$|q\rangle = \begin{pmatrix} a \\ b \end{pmatrix} = a \begin{pmatrix} 1 \\ 0 \end{pmatrix} + b \begin{pmatrix} 0 \\ 1 \end{pmatrix} = a |0\rangle + b |1\rangle$$

where $a, b \in \mathbb{C}$.

Products

The **Hermitian conjugate** of a quantum state is its complex transpose and is represented as follows: If

$$|q\rangle = \begin{pmatrix} a \\ b \end{pmatrix}$$

then

$$\langle q | \equiv |q\rangle^{\dagger} = \begin{pmatrix} a^* & b^* \end{pmatrix}$$

The inner product of two states $|\phi\rangle$ and $|\psi\rangle$ is given by $\langle\phi|\psi\rangle$.

Problem 1

Complete the following:

$$\langle 0|0\rangle = \begin{pmatrix} 1 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = 1$$

$$\langle 0|1\rangle = \begin{pmatrix} 1 & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = 0$$

$$\langle 1|0\rangle = \begin{pmatrix} 0 & 1 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} =$$

$$\langle 1|1\rangle =$$

As a general rule:

•
$$\langle a|b\rangle = 1$$
 if $a = b$

•
$$\langle a|b\rangle =$$

The **outer product** of two states ψ and ϕ is given by $|\psi\rangle |\phi\rangle$.

Problem 2

Complete the following:

$$|0\rangle \langle 0| = \begin{pmatrix} 1\\0 \end{pmatrix} \begin{pmatrix} 1&0 \end{pmatrix} = \begin{pmatrix} 1&0\\0&0 \end{pmatrix}$$
$$|0\rangle \langle 1| = \begin{pmatrix} 1\\0 \end{pmatrix} \begin{pmatrix} 0&1 \end{pmatrix} = \begin{pmatrix} 0&1\\0&0 \end{pmatrix}$$
$$|1\rangle \langle 0| = \begin{pmatrix} 0\\1 \end{pmatrix} \begin{pmatrix} 1&0 \end{pmatrix}$$
$$|1\rangle \langle 1| =$$

Problem 3

Finish the following mixed products by using the results of Problem 1:

Now finish complete the following rules:

- When multiplying the outer-product $|0\rangle \langle 0|$ by the state $|x\rangle$, the outer-product "checks" if x = 0 because
 - if x = 0, the state remains unaffected
 - if x = 1, the state is destroyed

• When multiplying the outer-product $|1\rangle \langle 1|$ by the state $|x\rangle$, the outer-product "checks" if x = 1 because

- if x = 0,

- if x = 1,
- The outer-product $|1\rangle \langle 0|$ can be though of as a "raising" operator because, when multiplied by the state $|x\rangle$
 - if x = 0, the outer-product "raises" the state from $|0\rangle$ to $|1\rangle$
 - if x = 1, the state is destroyed as 1 cannot be raised.
- The outer-product $|0\rangle \langle 1|$ can be though of as a "lowering" operator because, when multiplied by the state $|x\rangle$
 - if x = 0,
 - if x = 1,

The **tensor product** of two $n \times n$ matrices

$$A = \begin{pmatrix} a_{00} & a_{01} & \dots & a_{0n} \\ a_{10} & a_{11} & & \vdots \\ \vdots & & \ddots & \vdots \\ a_{n0} & \dots & \dots & a_{nn} \end{pmatrix}, B = \begin{pmatrix} b_{00} & b_{01} & \dots & b_{0n} \\ b_{10} & b_{11} & & \vdots \\ \vdots & & \ddots & \vdots \\ b_{n0} & \dots & \dots & b_{nn} \end{pmatrix}$$

is defined to be

$$A \otimes B = \begin{pmatrix} a_{00}B & a_{01}B & \dots & a_{0n}B \\ a_{10}B & a_{11}B & & \vdots \\ \vdots & & \ddots & \vdots \\ a_{n0}B & \dots & \dots & a_{nn}B \end{pmatrix}$$

Problem 4

The tensor product can be used to form a larger space of qubits. Using the notation $|ab\rangle = |a\rangle \otimes |b\rangle$ we can create a 2-qubit space by taking the tensor products of the single qubits. For example, if

$$|a\rangle = \begin{pmatrix} a_0\\a_1 \end{pmatrix} |b\rangle = \begin{pmatrix} b_0\\b_1 \end{pmatrix}$$

then

$$|ab\rangle = |a\rangle \otimes |b\rangle = \begin{pmatrix} a_0 \\ a_1 \end{pmatrix} \otimes \begin{pmatrix} b_0 \\ b_1 \end{pmatrix} = \begin{pmatrix} a_0 \begin{pmatrix} b_0 \\ b_1 \end{pmatrix} \\ a_1 \begin{pmatrix} b_0 \\ b_1 \end{pmatrix} = \begin{pmatrix} a_0b_0 \\ a_0b_1 \\ a_1b_0 \\ a_1b_1 \end{pmatrix}$$

Complete the following

$$|00\rangle = \begin{pmatrix} 1\\0 \end{pmatrix} \otimes \begin{pmatrix} 1\\0 \end{pmatrix} = \begin{pmatrix} 1\begin{pmatrix}0\\0\\0\\0 \end{pmatrix} \\ 0\begin{pmatrix}1\\0 \end{pmatrix} \\ 0\begin{pmatrix}0\\0 \end{pmatrix} \\ |01\rangle = \begin{pmatrix} 1\\0 \end{pmatrix} \otimes \begin{pmatrix} 0\\1 \end{pmatrix} = \begin{pmatrix} 1\begin{pmatrix}0\\1\\0\\0\\1 \end{pmatrix} \\ 0\begin{pmatrix}0\\1 \end{pmatrix} \\ 0\begin{pmatrix}0\\1 \end{pmatrix} \\ |10\rangle = \begin{pmatrix} 0\\1 \end{pmatrix} \otimes \begin{pmatrix} 1\\0 \end{pmatrix} = |11\rangle =$$

An important property of tensor products is the following

$$(A_1 \otimes A_2) (B_1 \otimes B_2) = A_1 B_1 \otimes A_2 B_2$$

Gates

A quantum algorithm can be written in terms of a **quantum circuit** which shows how qubits are manipulated by **quantum gates**, which are 2×2 , unitary matrices.

A gate U is represented in a quantum circuit as

$$|\psi\rangle$$
 — $U|\psi\rangle$

which means $|\psi\rangle \to U |\psi\rangle$.

Applying U to the state and then V

$$|\psi\rangle - U - V |\psi\rangle$$

is the same as $\left|\psi\right\rangle\rightarrow VU\left|\psi\right\rangle.$ Note the switched order.

The **Pauli gates** correspond to the Pauli matrices.

$$I = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$
$$X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$
$$Y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$
$$Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

To see how Pauli-gates are used to take expectation values, we'll write them in terms of outerproducts.

Problem 5

Complete the following:

Let's see how the X gate affects qubits.

$$X |0\rangle = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix} = |1\rangle$$
$$X |1\rangle = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \end{pmatrix} = |0\rangle$$

We've discovered that X "flips" the qubit that it acts on, from $|0\rangle$ to $|1\rangle$ and vice-verse. Therefore, we can write X in terms of outer-products as follows:

$$X = \left|1\right\rangle \left\langle 0\right| + \left|0\right\rangle \left\langle 1\right|$$

because then

$$X |0\rangle = (|1\rangle \langle 0| + |0\rangle \langle 1|) |0\rangle = |1\rangle \langle 0|0\rangle + |0\rangle \langle 1|0\rangle = |1\rangle \langle 1) + |0\rangle \langle 0) = |1\rangle$$
$$X |1\rangle = (|1\rangle \langle 0| + |0\rangle \langle 1|) |1\rangle = |1\rangle \langle 0|1\rangle + |0\rangle \langle 1|1\rangle = |1\rangle \langle 0) + |0\rangle \langle 1) = |0\rangle$$

Now let's look at the Y gate

$$Y |0\rangle = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ i \end{pmatrix} = i |1\rangle$$
$$Y |1\rangle = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} -i \\ 0 \end{pmatrix} =$$

Therefore, we can write Y in terms of outer-products as follows:

$$Y = i \left| 1 \right\rangle \left\langle 0 \right| - i \left| 0 \right\rangle \left\langle 1 \right|$$

because then

$$\begin{array}{l} Y\left|0\right\rangle = \left(i\left|1\right\rangle\left\langle 0\right| - i\left|0\right\rangle\left\langle 1\right|\right)\left|0\right\rangle = i\left|1\right\rangle\left\langle 0|0\right\rangle - i\left|0\right\rangle\left\langle 1|0\right\rangle = \\ Y\left|1\right\rangle = \left(i\left|1\right\rangle\left\langle 0\right| - i\left|0\right\rangle\left\langle 1\right|\right)\left|1\right\rangle = \end{array}$$

Finally, let's look at the Z gate

$$Z |0\rangle = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = Z |1\rangle =$$

Therefore, we can write Z in terms of outer-products as follows:

$$Z =$$

because then

$$Z |0\rangle = Z |1\rangle =$$

Bonus: I can be written in terms of outer-products as follows:

$$I =$$

We'll introduce two more gates, the Haddamard gate H and the phase gate S. They are defined as follows:

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1\\ 1 & -1 \end{pmatrix} \quad S = \begin{pmatrix} 1 & 0\\ 0 & i \end{pmatrix}$$

Problem 6

Prove that HZH = X

$$HZH = \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$
$$= \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix}$$
$$=$$

Prove that $SXS^{\dagger} = Y$

$$SXS^{\dagger} = \begin{pmatrix} 1 & 0 \\ 0 & i \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & i \end{pmatrix}^{\dagger}$$
$$= \begin{pmatrix} 1 & 0 \\ 0 & i \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -i \end{pmatrix}$$
$$=$$

Show that $Y = (HS^{\dagger})^{\dagger} Z (HS^{\dagger})$ using the previous two results and the identity $(AB)^{\dagger} = B^{\dagger} A^{\dagger}$.

Measurement

If a qubit is in the state

$$|q\rangle = \begin{pmatrix} a\\b \end{pmatrix}$$

then, when the qubit is measured

- the probability that $|0\rangle$ is measured is $P_{|\psi\rangle}(0) = |\langle 0|q\rangle|^2 = \left|\begin{pmatrix} 1 & 0 \end{pmatrix} \begin{pmatrix} a \\ b \end{pmatrix}\right|^2 = |a|^2$
- the probability that $|1\rangle$ is measured is $P_{|\psi\rangle}(1) = |\langle 1|q\rangle|^2 = \left|\begin{pmatrix} 0 & 1 \end{pmatrix} \begin{pmatrix} a \\ b \end{pmatrix}\right|^2 = |b|^2$

Expectation Value

Every observable (thing one can observe) has a corresponding operator O. If particle is in a state $|q\rangle$ then the expected value of the observable corresponding to the operator O is given by

$$\langle q|O|q\rangle$$

Expectation Values

The following will demonstrate how to find the expectation value of a tensor string of Pauli spin matrices.

Expectation Value of Z

The expectation value of Z in state $|\psi\rangle$ is give by

$$E_{\psi}(Z) = \langle \psi | Z | \psi \rangle$$

= $\langle \psi | (|0\rangle \langle 0| - |1\rangle \langle 1|) | \psi \rangle$
= $\langle \psi | 0 \rangle \langle 0 | \psi \rangle - \langle \psi | 1 \rangle \langle 1 | \psi \rangle$
= $|\langle 0 | \psi \rangle|^2 - |\langle 1 | \psi \rangle|^2$
= $P_{|\psi\rangle}(0) - P_{|\psi\rangle}(1)$

where $P_{|\psi\rangle}(x)$ is the probability that state $|\psi\rangle$ is measured to be $|x\rangle$.

So one creates the circuit

$$|\psi
angle$$
 —

which means prepare the state ψ and measure. One does this over and over, counting C_0 , the number of times one measures 0 and C_1 , the number of times one measures 1. From these numbers, on can estimate the probability of measuring 0 and 1 as

$$P_{|\psi\rangle}(0) = \frac{C_0}{C_0 + C_1}$$
$$P_{|\psi\rangle}(1) = \frac{C_1}{C_0 + C_1}$$

To find the expectation value of Z, one simply subtracts the probabilities as above:

$$P_{|\psi\rangle}(0) - P_{|\psi\rangle}(1) = \frac{C_0 - C_1}{C_0 + C_1}$$

Expectation Value of X

To find the expectation value of X, we are going to rotate our computation basis. Recall that X = HZH. Thus

$$E_{\psi}(X) = \langle \psi | X | \psi \rangle$$

$$= \langle \psi | HZH | \psi \rangle$$

$$= \langle H^{\dagger}\psi | Z | H\psi \rangle$$

$$= \langle H\psi | Z | H\psi \rangle$$

$$= E_{H\psi}(Z)$$

$$= P_{H|\psi\rangle}(0) - P_{H|\psi\rangle}(1)$$
(1)

So one creates the circuit



estimates the probabilities, and subtracts them, as above.

Problem 7

Complete the following:

Expectation Value of Y

To find the expectation value of Y, recall that $Y = SXS^{\dagger}$. Therefore $Y = (HS^{\dagger})^{\dagger} Z (HS^{\dagger})$. Thus

$$E_{\psi}(Y) = \langle \psi | Y | \psi \rangle$$

So one creates the circuit



estimates the probabilities, and subtracts them, as above.

To summarize

(Expectation value of:	Apply:	
J	Ζ	Ι	(2)
١	X	Н	
	$\langle Y \rangle$	HS^{\dagger}	

String of Pauli's

Let's work out how to take the expectation value of a tensor-product string of Pauli matrices. For example, let's work out how to take the expectation value of $X \otimes Z$.

$$\begin{aligned} E_{\psi}(X \otimes Z) &= \langle \psi | X \otimes Z | \psi \rangle \\ &= \langle \psi | (HZH) \otimes (IZI) | \psi \rangle \\ &= \langle \psi | (H \otimes I) (Z \otimes Z) (H \otimes I) | \psi \rangle \\ &= \langle (H \otimes I)^{\dagger} \psi | (Z \otimes Z) | (H \otimes I) \psi \rangle \\ &= \langle (H \otimes I) \psi | Z \otimes Z | (H \otimes I) \psi \rangle \\ &= \langle \psi' | (|0\rangle \langle 0| - |1\rangle \langle 1|) \otimes (|0\rangle \langle 0| - |1\rangle \langle 1|) | \psi' \rangle \\ &= \langle \psi' | (|00\rangle \langle 00| - |01\rangle \langle 01| - |10\rangle \langle 10| - |11\rangle \langle 11|) | \psi' \rangle \\ &= \langle \psi' | 00\rangle \langle 00| \psi' \rangle - \langle \psi' | 01\rangle \langle 01| \psi' \rangle \\ &- \langle \psi' | 10\rangle \langle 10| \psi' \rangle + \langle \psi' | 11\rangle \langle 11| \psi' \rangle \\ &= |\langle 00| \psi' \rangle|^2 - |\langle 01| \psi' \rangle|^2 - |\langle 10| \psi' \rangle|^2 + |\langle 11| \psi' \rangle|^2 \\ &= P_{\psi'}(00) - P_{\psi'}(01) - P_{\psi'}(10) + P_{\psi'}(11) \end{aligned}$$

where I've defined $\psi' = (H \otimes I)\psi$.

Problem 8

Work out the following expectation value:

 $E_{\psi}(I \otimes X) = \langle \psi | I \otimes X | \psi \rangle$

This can be extended to tensor-strings of arbitrary length. We can also extend this to sums of tensor-strings because expectation value is linear. For example:

$$\langle \psi | (X+Z) | \psi \rangle = \langle \psi | X | \psi \rangle + \langle \psi | Z | \psi \rangle$$

The power of this is that, if one can write a Hamiltonian H in terms of a linear combination of tensor-strings of Pauli matrices, one can use a quantum computer to estimate the expectation values of H (the energy) of the system for a given state $|\psi\rangle$.

We are now equipped to learn about the variation quantum eigensolver.